

## Conductance of a Perfect Thin Film with Diffuse Surface Scattering

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The conductance of thin films with diffusive surface scattering was solved semiclassically by Fuchs and Sondheimer. However, when the intrinsic electron mean free path is very large or infinite their conductance diverges. In this Letter a simple diffraction picture is presented. It yields a conductance which corresponds to a limiting mean free path of  $a^2/\lambda_F$ , where  $a$  is the film thickness.

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The conductivity of a thin film generally decreases considerably when the film thickness is reduced. The main reason is that the surfaces of the film scatter the conduction electrons partially or completely diffusively and reduce their effective mean free path. This phenomenon was first treated by Fuchs [1] and later generalized by Sondheimer [2]. The result of their semiclassical calculation for the conductance  $G$  of a square shaped thin film (width equal to length) of thickness  $a$  with completely diffusive scattering at the surfaces is

$$G_{\square} = \frac{ne^2 a l_0}{m v_F} \left[ 1 - \frac{3l_0}{2a} \int_0^1 (1-u^2) u \left\{ 1 - \exp\left(-\frac{a}{l_0 u}\right) \right\} du \right], \quad (1)$$

where  $n$  is the density,  $l_0$  the bulk mean free path, and  $v_F$  the Fermi velocity of the electrons.  $u = \sin\alpha$ , where  $\alpha$  is the angle between the velocity vector of the electrons and the film plane. Chambers [3] interpreted these results by introducing an  $\mathbf{r}$ - and  $\mathbf{k}$ -dependent vector mean free path  $\mathbf{l}_{\mathbf{r},\mathbf{k}} = \int_{-t_0}^0 \mathbf{v}(t') e^{t'/\tau} dt'$ . The vector mean free path in a film is reduced by the diffusive scattering at the surface (at the earlier time  $-t_0$ ) and depends therefore on the position in the film and the  $z$  component of the Fermi velocity. While the conduction electron travels along the vector mean free path it accumulates momentum and energy. The energy is equal to  $(-e)\mathbf{E} \cdot \mathbf{l}_{\mathbf{r},\mathbf{k}}$  and yields a correction  $f_1(\mathbf{r}, \mathbf{k}) = (-\frac{df_0}{de})(-e)\mathbf{E} \cdot \mathbf{l}_{\mathbf{r},\mathbf{k}}$  to the distribution function  $f(\mathbf{r}, \mathbf{k}) = f_0 + f_1(\mathbf{r}, \mathbf{k})$ , where  $f_0$  is the equilibrium Fermi function.

The Fuchs-Sondheimer conductance of a thin film diverges when the intrinsic mean free path becomes infinite. For a very large intrinsic mean free path the conductance of a thin film becomes [2]

$$G_{\square} = \frac{ne^2}{m v_F} a * \frac{3a}{4} \ln \frac{l_0}{a} = \left(\frac{a}{\lambda_F}\right)^2 \frac{e^2}{\hbar} \ln \frac{l_0}{a}, \quad (2)$$

where  $\lambda_F = 2\pi/k_F$  is the Fermi wavelength.

The divergence is due to the electrons which travel almost parallel to the surface, so that their effective mean free path becomes very large. An electron with  $v_z = v_F \sin\alpha$  travels the distance  $av_F/v_z = a/\sin\alpha \approx a/\alpha$  from one surface to the opposite one if there are no internal collisions. For a given  $v_z$  the distance traveled since the last

collision increases linearly with the distance from the surface of the last collision from 0 to  $av_F/v_z$ . Therefore the average mean free path for electrons with a given  $v_z$  or a given angle  $\alpha$  is half the distance between two surface collisions  $l_{ef} = a/(2\alpha)$ . This  $l_{ef}$  diverges for  $\alpha \rightarrow 0$ .

It has been pointed out in the past that this divergence is a classical effect which should disappear if one treats the problem quantum theoretically. Tesanovic *et al.* [4] performed such a calculation. They considered a thin film whose thickness is modulated (modeling a finite surface roughness) and transformed the surface modulation into a scattering potential. In their evaluation they expanded the wave function in the basis of standing waves in the  $z$  direction perpendicular to the film plane,  $\psi \propto \sin(k_z z)$ . This yields a quantization of  $k_z = n_z \pi/a$ ,  $n_z = 1, 2, \dots$ . Comparison with experimental data yielded a reasonable modulation amplitude for the film thickness.

For a completely diffusive scattering at the surface the average (vector) mean free path in the  $z$  direction is equal to half the film thickness  $a/2$ . This yields an uncertainty of  $\Delta k_z$  which is about  $2/a$ , i.e., of the order of the separation of the  $k_z$  values. Therefore this approach might not be well suited for a surface with completely diffusive scattering.

In this Letter we give (i) a semiquantitative evaluation of the scattering in the Feynman path integral approach and (ii) treat the problem as an electron diffraction by the film boundaries. Both treatments yield essentially the same quantitative results that the diffusive surface scattering acts as if the conduction electrons experience an effective mean free path of  $l_0 = a^2/(2\lambda_F)$  ( $\lambda_F = 2\pi/k_F =$  Fermi wavelength).

The Kubo formalism provides a tool to calculate the conductance for an arbitrary geometry. In the real space representation the Kubo formula uses the Green function  $G_{\omega}(\mathbf{r}, \mathbf{r}')$  for the propagator from  $\mathbf{r}$  to  $\mathbf{r}'$ . If one uses the following form of the real space Green function,

$$G_{\omega}(\mathbf{r} - \mathbf{r}') = \frac{m}{2\pi\hbar} \frac{1}{|\mathbf{r} - \mathbf{r}'|} \times \left[ \exp\left(ik_F |\mathbf{r} - \mathbf{r}'| \text{sgn}(\omega) - \frac{|\omega|}{v_F} |\mathbf{r} - \mathbf{r}'|\right) \right], \quad (3)$$

then one can reproduce the vector mean free path result for the conductance (which is identical with the Fuchs-Sondheimer formula).

In a thin film with diffuse surface scattering the Green function can no longer be expressed by Eq. (3). The reason can be understood in terms of Feynman's path integral. A (quantum theoretical) electron does not propagate along a straight line with constant velocity  $v_F$  from  $\mathbf{r}$  to  $\mathbf{r}'$  but uses all possible paths from  $\mathbf{r}$  to  $\mathbf{r}'$ . Of course, most of these paths are not important because their contribution averages to zero due to their phase factor  $\exp[\frac{i}{\hbar} \int_0^T L(\mathbf{r}, \mathbf{r}^\circ) dt]$ , [ $L(\mathbf{r}, \dot{\mathbf{r}}) = \frac{m}{2}(\dot{\mathbf{r}})^2$  is the Lagrangian]. However, paths which are sufficiently close to a straight line contribute to the amplitude since one has to sum over all paths.

In Fig. 1 a thin film is drawn with the position  $\mathbf{r} = (-L/2, 0, 0)$  in the center of the film and the position  $\mathbf{r}' = (L/2, 0, 0)$  at a distance  $L$  from  $\mathbf{r}$ . The film plane is perpendicular to the  $z$  direction and the electric field is applied in the  $x$  direction. The traveling time for an electron with Fermi velocity is  $T = L/v_F$ , and the phase of the direct path is  $\phi_0 = (m/2)v_F^2 T/\hbar$ . If we consider a very simple alternative path in which the electron moves from  $\mathbf{r}$  to  $\mathbf{r}'' = (0, 0, z)$  in the time interval  $T/2$  and then from  $\mathbf{r}''$  to  $\mathbf{r}'$  its average velocity is  $v' = v_F \sqrt{1 + (z/2L)^2}$ . This yields for the phase of this path  $\phi = (m/2)v_F^2 [1 + (z/2L)^2] T/\hbar = \phi_0 + k_F z^2/(8L)$ . As long as  $\Delta\phi = k_F z^2/(8L) < \pi/2$  this path and similar ones contribute to the amplitude at  $\mathbf{r}'$ .

In Fig. 1 we have chosen  $L = k_F a^2/(4\pi)$  and drawn an ellipse with focal points at  $(\pm L/2, 0, 0)$  and the short semi-axis equal to the film thickness  $a$ . All paths  $\mathbf{r} \rightarrow \mathbf{r}'' \rightarrow \mathbf{r}'$  where  $\mathbf{r}''$  lies within the ellipse have a phase difference which is less than  $\pi/2$  with respect to the direct path; all paths where  $\mathbf{r}''$  lies outside the ellipses have a phase difference larger than  $\pi/2$ . For this situation a large fraction of the constructive paths intersects the film boundary and is eliminated because of the diffuse scattering at the surface. We can roughly say that the fraction of electrons which reach the point  $\mathbf{r}'$  without phase randomization is reduced by a factor of 2. That means that the mean free path of the electrons parallel to the film is of the order of  $k_F a^2/(4\pi)$ . So their contribution to the conductance is finite. One obtains roughly the same reduction when the  $z$  components of  $\mathbf{r}$  and  $\mathbf{r}'$  are varied between  $\pm a/2$ . This means that all paths within the angle  $\sin\alpha = a/L =$

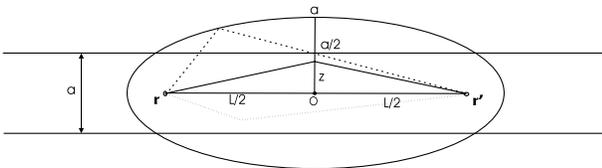


FIG. 1. Different Feynman paths from  $\mathbf{r}$  to  $\mathbf{r}'$  in a thin film via an intermediate point  $\mathbf{r}''$ . All these paths inside the ellipse are in phase with the direct path. The diffusive scattering at the film surfaces eliminates many of these paths.

$4\pi/(ak_F) = 2\lambda_F/a$  have a finite mean free path of about  $a^2/(2\lambda_F)$ .

The quantitative evaluation of the Feynman path integral is a rather complex task, in particular, the normalization of the sum or integral over the different paths. Therefore we turn to a different and quite simple approach which I call the diffraction model.

If we cut the film at a plane parallel to the  $y$ - $z$  plane, then the cross section of the film represents a slit with the width  $a$ . This is shown in Fig. 2. Through this slit the conduction electron waves propagate with the wavelength  $\lambda_F = \frac{2\pi}{k_F}$ . Therefore the electron waves experience diffraction. The result of the diffraction is that the wave spreads out. The angle for the first diffraction minimum is

$$\sin\alpha_d = \frac{\lambda_F}{a} = \frac{2\pi}{k_F a}.$$

Now we consider an electron which travels parallel to the film. Then the slit experiment tells us that at the distance  $dx = v_F dt$  after the slit the beam of width  $a$  has broadened to  $(a + 2\alpha_d dx) = (a + 2\alpha_d v_F dt)$ . Since the fraction of the wave outside of  $a$  is scattered at the surface the remaining intensity  $I$  of the electron within the film reduces to  $(I + dI)$ ,  $dI < 0$ ,

$$I + dI = I \frac{a}{a + 2\alpha_d dx} \approx I \left(1 - \frac{2\alpha_d}{a} dx\right),$$

$$dI = -I \frac{2\alpha_d}{a} dx = -I \frac{2\lambda_F}{a^2} dx = -I \frac{2\lambda_F v_F}{a^2} dt,$$

$$I = I_0 \exp\left(-\frac{2\alpha_d x}{a}\right) = I_0 \exp\left(-\frac{2\alpha_d v_F t}{a}\right).$$

As a consequence the mean free path of the electron which travels parallel to the film is  $l_d = \frac{a}{2\alpha_d} = \frac{a^2}{2\lambda_F}$  or the relaxation time is  $\tau = \frac{a}{2\alpha_d v_F} = \frac{a^2 m}{4\pi\hbar}$ . Since the diffraction of the electron wave occurs on both film surfaces within an angle  $\alpha_d$  all electrons which travel almost parallel to the  $x$ - $y$  plane within the angle  $\alpha_d$  have the same mean free

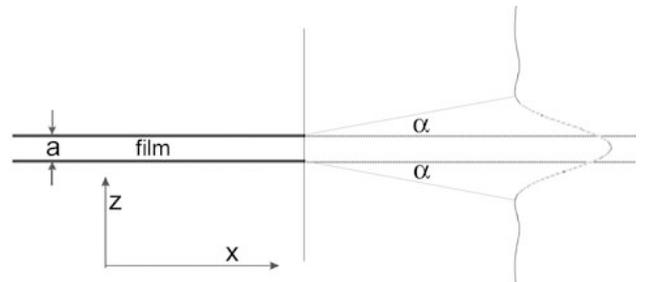


FIG. 2. The film is cut perpendicular to show the spread of the electron wave function due to diffraction by a single slit. The curve on the right gives the single slit diffraction intensity. The part of the electron wave outside of the film (dashed lines) will be diffusively scattered. Therefore the electron wave is damped.

path  $l_d = a^2/(2\lambda_F)$ . Actually  $l_d$  is identical with half the distance it takes an electron to travel from one surface to the opposite one at the angle  $\alpha_d$ . For larger angles  $\alpha$  between the wave vector and the film plane the diffraction has little effect, even a classical electron hits the surface and the mean free path (averaged over the position in the film) is  $a/(2 \sin \alpha)$  which is shorter than  $a/(2\alpha_d)$ .

We obtain the quantum theoretical conductance of a thin film with diffuse surface scattering but an infinite intrinsic mean free path by dividing the integration over  $du = d(\sin \alpha)$  into two regions; for  $|\alpha| > |\alpha_d|$  we use an infinite mean free path and for  $|\alpha| < |\alpha_d|$  the finite mean free path of  $a^2/(2\lambda_F)$ . We find for  $G_\infty$ ,

$$G_\infty = \frac{e^2}{\hbar} \left\{ \left[ \left( \frac{a}{\lambda_F} \right)^2 \left( \ln \frac{a}{\lambda_F} + \frac{1}{2} \right) + \frac{1}{6} \right] \right\}. \quad (4)$$

This result has logarithmic accuracy only. Therefore we ignore the nonlogarithmic terms and obtain as our final result

$$G_\infty \cong \frac{e^2}{\hbar} \left( \frac{a}{\lambda_F} \right)^2 \left( \ln \frac{a}{\lambda_F} \right). \quad (5)$$

The conductance depends only on the ratio  $a/\lambda_F$ .

Obviously the conductance does not diverge in contrast to the semiclassical solution by Fuchs and Sondheimer. Because of the wave character of the electrons the diffuse surface scattering introduces an effective mean free path. One can incorporate this result in the semiclassical solution by replacing the intrinsic mean free path in Eq. (2) by  $a^2/\lambda_F$ .

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- [1] K. Fuchs, Proc. Cambridge Philos. Soc. **34**, 100 (1938).
  - [2] E. H. Sondheimer, Adv. Phys. **50**, 499 (2001).
  - [3] R. G. Chambers, in *The Physics Of Metals*, edited by J. M. Ziman (Cambridge University Press, London, 1969), p. 175.
  - [4] Z. Tesanovic, M. V. Jaric, and S. Maekawa, Phys. Rev. Lett. **57**, 2760 (1986).